

Assignment 2 — Solutions [*Revision* : 1.2]

Question 1

To obtain the desired density-radius relation, we need to derive an expression for the luminosity in terms of the density and radius, and then combine this with the empirical $L \propto R^3$ luminosity-radius relation.

The stellar luminosity is found by integrating the equation of energy conservation, viz

$$L = \int_0^R \frac{dL_r}{dr} dr = \int_0^R 4\pi r^2 \rho \epsilon dr. \quad (1)$$

The nuclear energy release rate depends on the density ρ (assumed constant throughout the star) and the temperature T , via

$$\epsilon = \epsilon_0 \rho T^4 \quad (2)$$

(where ϵ_0 is a constant). Thus, to calculate the stellar luminosity, we're going to need an expression for the temperature as a function of radius.

This expression can be found by first solving the equation of hydrostatic equilibrium,

$$\frac{dP}{dr} = -\frac{GM_r}{r^2} \rho. \quad (3)$$

For a constant-density star, the mass coordinate M_r is simply the mass of a uniform sphere with radius r ,

$$M_r = \frac{4\pi r^3 \rho}{3}. \quad (4)$$

Hence,

$$\frac{dP}{dr} = -\frac{4\pi G r \rho^2}{3}, \quad (5)$$

and, adopting the surface boundary condition $P(R) = 0$, the pressure distribution throughout the star is

$$P(r) = \frac{2\pi G \rho^2}{3} (R^2 - r^2). \quad (6)$$

Using the ideal gas law, this is readily transformed into a temperature distribution

$$T(r) = \frac{2\pi G \rho \mu u}{3k} (R^2 - r^2). \quad (7)$$

Combining the above expression with that for ϵ , the luminosity integral (1) can be evaluated as

$$L = \frac{8192 G^4 \pi^5 R^{11} u^4 \epsilon_0 \mu^4 \rho^6}{280665 k^4}, \quad (8)$$

or more simply,

$$L \propto R^{11} \rho^6. \quad (9)$$

With the $L \propto R^3$ relation, this becomes

$$R^3 \propto R^{11} \rho^6, \quad (10)$$

and so

$$\rho \propto R^{-4/3}. \quad (11)$$

Hence, the exponent in the radius-density relation is $\chi = -4/3$, and — because this exponent is negative — we find that bigger/brighter stars are less dense than smaller/dimmer stars. Voila!