Homework Assignment 6 — Solutions

• Q8.1

Room temperature is 294 K, so $kT = 4.06 \times 10^{-14} \text{ J} = 0.0253 \text{ eV} \approx 1/40 \text{ eV}$. kT = 1 eV when T = 11,600 K, and kT = 13.6 eV when T = 158,000 K.

• Q8.5

The Boltzmann equation gives us

$$\frac{n_2}{n_1} = \frac{g_2}{g_1} e^{-(E_2 - E_1)/kT} = 4 e^{-10.2 \,\mathrm{eV}/kT},$$

so we have $\frac{n_2}{n_1} = 1\%$ when $e^{-10.2 \text{ eV}/kT} = 0.0025$. Taking the logarithm of both sides gives -10.2 eV/kT = -5.99, and solving for T gives $T = 1.97 \times 10^4 \text{ K}$. The same procedure for $\frac{n_2}{n_1} = 10\%$ gives $T = 3.21 \times 10^4 \text{ K}$.

- Q8.6
 - (a). From the Boltzmann equation,

$$\frac{n_3}{n_1} = \frac{g_3}{g_1} e^{-(E_3 - E_1)/kT} = 9 e^{-12.1 \,\mathrm{eV}/kT},$$

which gives $\frac{n_3}{n_1} = 1$ when $T = 6.38 \times 10^4 \, \text{K}.$

- (b). When T = 85,400 K, $\frac{N_3}{N_1} = 9 e^{-1.64} = 1.74$, so if there are $N_1 = N$ atoms in the n = 1 state then there are $N_3 = 1.74 N$ atoms in the second excited state (n = 3).
- (c). As $T \to \infty$, the exponential factor in the Boltzmann equation becomes 1 for every value of n, so the Boltzmann equation predicts that the distribution of electrons mirrors the values of the degeneracies: level n has a number of electrons proportional to n^2 .

In reality, as $T \to \infty$ all atoms ionize, and all the electrons are free electrons.

• Q8.7

We have $Z_{\rm I} = g_1 + g_2 e^{-10.2 \,{\rm eV}/kT} + g_3 e^{-12.1 \,{\rm eV}/kT}$, with $g_1 = 2(1)^2 = 2$, $g_2 = 2(2)^2 = 8$, and $g_3 = 2(3)^2 = 18$. At $T = 10,000 \,{\rm K}$, this gives $Z_{\rm I} = 2 + 5.81 \times 10^{-5} + 1.46 \times 10^{-5}$. We can see that the second and third terms are much smaller than $g_1 = 2$, so we have $Z_{\rm I} \approx 2$.

- Q8.9
 - (a). The Saha equation is

$$\frac{N_{\rm II}}{N_{\rm I}} = \frac{2Z_{\rm II}}{n_e Z_{\rm I}} \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT}.$$

We also have charge conservation

$$n_e V = N_{II}$$

and conservation of nucleons

$$N_t = \rho V / (m_p + m_e) \approx \rho V / m_p,$$

where $N_t = N_{\rm I} + N_{\rm II}$ is the total number of hydrogen nuclei (whether in neutral hydrogen or ionized hydrogen).

The first step is to replace the partition function $Z_{\rm I}$ with the ground state degeneracy $g_1 = 2$ —recall from question 8.7 that this is a good approximation at 10,000 K, and in fact it is a good

approximation for $kT \ll 10.2 \text{ eV}$, or $T \ll 120,000 \text{ K}$. The partition function Z_{II} of the proton is 1. We can now write the Saha equation (for this pure hydrogen case, and assuming $T \ll 120,000 \text{ K}$) as

$$\frac{N_{\rm II}}{N_{\rm I}} = \frac{1}{n_e} \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT} = \frac{V}{N_{\rm II}} \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT},$$

where we use charge conservation to get rid of n_e in the final expression. We also have $N_t = N_{\rm I} + N_{\rm II} = N_{\rm I} (1 + N_{\rm II}/N_{\rm I})$; from the Saha equation, we get $N_t = N_{\rm I} (1 + x/N_{\rm II})$, where $x = V \left(\frac{2\pi m_e kT}{\hbar^2}\right)^{3/2} e^{-\chi_{\rm I}/kT}$. We can rearrange this as $N_{\rm I} = N_t / (1 + x/N_{\rm II})$. Putting this in the Saha equation gives

$$\frac{N_{\rm II}}{N_t / \left(1 + x / N_{\rm II}\right)} = \frac{x}{N_{\rm II}}$$

or

$$\frac{N_{\mathrm{II}}^2}{N_t}\left(1+x/N_{\mathrm{II}}\right)=x.$$

This gives

$$\frac{N_{\mathrm{II}}^2}{N_t} + x \frac{N_{\mathrm{II}}}{N_t} - x = 0.$$

Dividing through by N_t , we have

$$\left(\frac{N_{\rm II}}{N_t}\right)^2 + \frac{x}{N_t} \left(\frac{N_{\rm II}}{N_t}\right) - \frac{x}{N_t} = 0. \tag{1}$$

Now we use the conservation of nucleons, $N_t \approx \rho V/m_p$ and our expression for $x, x = V \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_1/kT}$, to get

$$\frac{x}{N_t} = \frac{V\left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT}}{\rho V/m_p} = \left(\frac{m_p}{\rho}\right) \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT}$$

which we plug in to equation 1 to get

$$\left(\frac{N_{\rm II}}{N_t}\right)^2 + \left(\frac{N_{\rm II}}{N_t}\right) \left(\frac{m_p}{\rho}\right) \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT} - \left(\frac{m_p}{\rho}\right) \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm I}/kT} = 0.$$

(b). Taking $\rho = 10^{-6} \text{ kg m}^{-3}$, we have

$$\left(\frac{N_{\rm II}}{N_t}\right)^2 + 4.04 \left(T/1\,{\rm K}\right)^{3/2} e^{-158,000\,{\rm K}/T} \left(\frac{N_{\rm II}}{N_t}\right) - 4.04 \left(T/1\,{\rm K}\right)^{3/2} e^{-158,000\,{\rm K}/T} = 0$$

Defining $b \equiv 4.04 (T/1 \text{ K})^{3/2} e^{-158,000 \text{ K}/T}$, we have the quadratic equation $y^2 + by - b = 0$, where y is the ionization fraction N_{II}/N_t . When $b \gg 1$, the quadratic term becomes negligible (remember, y must be between 0 and 1, so y^2 is also between 0 and 1), and the equation reduces to $by - b \approx 0$, which has solution $y \approx 1$. When $b \ll 1$, the linear and constant terms become negligible and we have $y^2 \approx 0$, so $y \approx 0$. Plugging in T = 5000 K gives $b = 2.7 \times 10^{-8} \ll 1$, and plugging in T = 25,000 K gives $b = 2.9 \times 10^4 \gg 1$, so we expect that our graph will go from $y = N_{\text{II}}/N_t \approx 0$ at the low temperature end to $N_{\text{II}}/N_t \approx 1$ at the high temperature end, with a transition at temperatures where $b \approx 1$ (in fact, $N_{\text{II}}/N_t = 0.5$ when b = 0.5). Looking at the plot, we see that these predictions are correct—the transition from fully neutral to fully ionized happens at a temperature of about 9900 K, at which point $b \approx 0.5$.

• Q8.13

From Example 8.1.5, we have T = 5777 K and $P_e = 1.5 \text{ N m}^{-2}$ in the solar photosphere, and $Z_{\text{II}} = 2.30$ for calcium. We are also given $\chi_{\text{II}} = 11.9 \text{ eV}$ and $Z_{\text{III}} = 1$. With this, we can write the Saha equation for the ratio of doubly-ionized calcium to singly-ionized calcium:

$$\frac{N_{\rm III}}{N_{\rm II}} = \frac{2kTZ_{\rm III}}{P_e Z_{\rm II}} \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_{\rm II}/kT} = 0.002.$$

Very little calcium is doubly ionized in the Sun's photosphere; instead, the calcium is almost all in the singly-ionized state—Example 8.1.5 shows that the ratio of singly-ionized calcium to neutral calcium is very high, and that almost all of the singly-ionized calcium is in the ground state. This means that almost all of the calcium in the Sun's photosphere can contribute to the formation of the H and K lines.

• Q8.14

The Saha equation is

$$\frac{N_{i+1}}{N_i} = \frac{2Z_{i+1}}{n_e Z_i} \left(\frac{2\pi m_e kT}{h^2}\right)^{3/2} e^{-\chi_i/kT}$$

We can see that the ratio N_{i+1}/N_i increases as the temperature T increases, and that it decreases as the electron density n_e increases. If we consider a main-sequence star and a giant star of the same spectral type, the fact that they have the same spectral type means that they should have (nearly) the same ratio N_{i+1}/N_i for all species so that they form the spectral lines with the same relative strengths. The fact that the giant star has a lower atmospheric density means that it has a lower value of n_e than the main-sequence star, so in order to maintain the same value of N_{i+1}/N_i as the main-sequence star the giant star must also have a lower value of T than the main-sequence star.

• Q8.16

Fomalhaut has V = 1.19; from the H-R diagram, it also has $M_V \approx 2$. We can solve for the distance to Fomalhaut using $V - M_V = 5 \log_{10} \left(\frac{d}{10 \text{ pc}}\right)$, which gives $d \approx 7 \text{ pc}$.

