

## 34 — Degeneracy [*Revision* : 1.1]

- Breakdown of ideal gas law

- So far, assumed gas pressure follows ideal equation of state:

$$P_{\text{gas}} = \frac{\rho kT}{\mu m_{\text{H}}}$$

- However, at sufficiently high densities, and sufficiently low temperatures, ideal assumption begins to break down:

- \* Number density of particles  $n$  is large
- \* Typical momentum of particles is small
- \* Particles concentrated in small volume of phase (momentum/position) space
- \* But there's a limit to how tightly particles can be packed in phase space: **Pauli exclusion principle**
- \* In limit where exclusion principle is important, gas is **degenerate**; different equation of state

- Non-relativistic degenerate gas

- Consider Maxwell-Boltzmann velocity distribution function for (non-degenerate) particles of mass  $m$  at temperature  $T$ :

$$n_v dv = n \left( \frac{m}{2\pi kT} \right)^{3/2} e^{-mv^2/2kT} 4\pi v^2 dv$$

(number of particles per unit volume with velocities in interval  $[v, v + dv]$ )

- Can also be written as momentum distribution function:

$$n_p dp = n \left( \frac{1}{2\pi mkT} \right)^{3/2} e^{-p^2/2mkT} 4\pi p^2 dp$$

- In limit  $T \rightarrow 0$ , particles concentrated around  $p = 0$ ; all try to have lowest possible momentum & energy  $E(p)$ , which is zero
- However, this can run foul of Pauli exclusion principle, since all particles in same momentum state
- In fact, lowest momentum/energy state has all particles filling up available cells of phase space, up to some maximum momentum  $p_{\text{F}}$  (the **Fermi momentum**)

$$n_p dp = \begin{cases} \frac{4\pi p^2 dp}{h^3/2} & p < p_{\text{F}} \\ 0 & p > p_{\text{F}} \end{cases}$$

Numerator: volume of shell in phase space with momentum in interval  $[p, p + dp]$ ; denominator: volume in phase space occupied by each distinct quantum state (spin gives factor of 2)

- Fermi momentum set by requirement

$$\int_0^{p_{\text{F}}} \frac{4\pi p^2 dp}{h^3/2} = n,$$

so that

$$p_{\text{F}} \sim n^{1/3}$$

(henceforth, drop all unimportant factors in expressions, for simplicity)

- To obtain equation of state, recall that pressure scales as

$$P_{\text{gas}} \sim u$$

where  $u$  is kinetic energy density

- To find  $u$ , integrate over all occupied states

$$u \sim \int_0^{p_F} E(p) n_p dp$$

If gas is non-relativistic,  $E(p) = p^2/2m$ ; hence,

$$u \sim \int_0^{p_F} p^2 p^2 dp \sim p_F^5$$

and

$$P_{\text{gas}} \sim p_F^5 \sim n^{5/3}$$

- With  $\rho \sim n$ , final result:

$$P_{\text{gas}} \sim \rho^{5/3}$$

- Discussion:

- \* This is **degeneracy pressure**; comes from exclusion principle rather than thermal motions
- \* Independent of temperature
- \* Polytropic — a fully degenerate star is a polytrope
- \* Strictly applies at zero temperature, but becomes good approximation when average particle momenta are below  $p_F$
- \* This limit equivalent to inequality

$$kT \lesssim \frac{p_F^2}{2m} \equiv E_F$$

( $E_F$  is **Fermi energy**). Note mass dependence; lighter particles become degenerate first

- \* Often see this inequality written as

$$\frac{T}{\rho^{2/3}} < \mathcal{D}$$

where  $\mathcal{D}$  depends on  $m$

- \* **Electron degeneracy** important in cores of low-mass stars, and in white dwarfs
- \* **Neutron degeneracy** important in neutron stars (no electrons left!)
- \* Degeneracy responsible for **helium flash**: when helium ignites, temperature increases (particles move faster), but pressure does not increase (because degeneracy condition above still holds); no 'safety valve' where star expands to cool off reactions, thus runaway burning

- Relativistic degenerate gas

- Similar to derivation above, but must use Einstein energy-momentum relation

$$E^2 = p^2 c^2 + m_0^2 c^4$$

( $m_0$  is rest mass)

– In relativistic limit  $E \gg m_0 c^2$ ,  $E \sim p$

– Hence,

$$u \sim \int_0^{p_F} p p^2 dp \sim p_F^4$$

and

$$P_{\text{gas}} \sim \rho^{4/3}$$

– Electrons become relativistic in white dwarf stars as they approach **Chandrasekhar limit** (limiting mass, above which there is collapse to neutron star)