9 — Describing a Radiation Field [*Revision* : 1.3]

• Specific Intensity

- To fully describe a radiation field, we need to specify how much energy
 - $\ast\,$ At each point in space
 - * At each point in time
 - * In each direction
 - $\ast\,$ At each wavelength
- All this information encapsulated in specific intensity I_{λ}
- For an infinitessimal area element dA, the amount of energy passing through dA at an angle θ to the normal, within the wavelength interval $(\lambda, \lambda + d\lambda)$, within the truncated cone with solid angle $d\Omega$, and within the time interval dt is:

$$E_{\lambda} \, \mathrm{d}\lambda = I_{\lambda} \, \mathrm{d}A \cos\theta \, \mathrm{d}\lambda \, \mathrm{d}\Omega \, \mathrm{d}t$$

- Aside: **solid angle** is 3-dimensional analog to planar angle.
 - * For circle of radius r, segment with angle ϕ (in radians) has arc length

 $\mathrm{d}s = r\,\phi$

* For sphere of radius r, cone with solid angle $d\Omega$ (in **steradians**) has base area

 $\mathrm{d} A = r^2 \mathrm{d} \Omega$

- * Full circle has total angle $2\pi\,\mathrm{rad};$ full sphere has $4\pi\,\mathrm{sterad}$
- * In spherical-polar coordinates, solid angle differential d Ω can be written in terms of $\theta,\,\phi$ differentials

$$\mathrm{d}\Omega = \sin\theta\,\mathrm{d}\theta\,\mathrm{d}\phi$$

- Mean intensity & energy density
 - At each point in space, define mean intensity by averaging over all solid angles

$$\langle I_{\lambda} \rangle = \frac{1}{4\pi} \int I_{\lambda} d\Omega = \frac{1}{4\pi} \int_{0}^{2\pi} \int_{0}^{\pi} I_{\lambda} \sin \theta \, \mathrm{d}\theta \, \mathrm{d}\phi$$

- For axisymmetric radiation field $(I_{\lambda} \text{ not depending on } \phi)$,

$$\langle I_{\lambda} \rangle = \frac{1}{2} \int_0^{\pi} I_{\lambda} \sin \theta \mathrm{d}\theta.$$

This often written

$$\langle I_{\lambda} \rangle = \frac{1}{2} \int_{-1}^{1} I_{\lambda} \mathrm{d}\mu$$

where $\mu \equiv \cos \theta$ (and $d\mu = \sin \theta d\theta$)

- For isotropic radiation field, $\langle I_{\lambda} \rangle = I_{\lambda}$
- Mean intensity is related to the **specific energy density** u_{λ} via

$$u_{\lambda} \mathrm{d}\lambda = \frac{4\pi}{c} \langle I_{\lambda} \rangle \mathrm{d}\lambda$$

- For blackbody (recall from notes 3):

$$u_{\lambda} \mathrm{d}\lambda = \frac{8\pi hc}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1} \,\mathrm{d}\lambda$$

Since BB radiation field is isotropic,

$$I_{\lambda} d\lambda = \langle I_{\lambda} \rangle d\lambda = \frac{c}{4\pi} u_{\lambda} d\lambda = \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1} d\lambda$$

This last equation defines **Planck function** — special name for specific intensity of BB radiation field:

$$B_{\lambda}(T) \equiv \frac{2hc^2}{\lambda^5} \frac{1}{e^{hc/\lambda kT} - 1}$$

(B for blackbody!)

- Total (bolometric) energy density for BB found by integrating over all wavelengths:

$$u = \int_0^\infty u_\lambda \mathrm{d}\lambda = \int_0^\infty B_\lambda \mathrm{d}\lambda = aT^4$$

where $a = 4\sigma/c$ is radiation constant

• Flux

- Total energy passing through surface, per second, per unit area, per unit wavelength, in all directions, defines **specific flux** F_{λ} (aka the monochromatic flux)
- To find F_{λ} , integrate equation for energy passing through dA over all solid angles:

$$F_{\lambda} = \frac{\int E_{\lambda}}{\mathrm{d}A\mathrm{d}t\mathrm{d}\lambda} = \int I_{\lambda}\cos\theta\mathrm{d}\Omega$$

Using definition of differential solid angle:

$$F_{\lambda} = \int_{0}^{2\pi} \int_{0} I_{\lambda} \cos \theta \sin \theta \mathrm{d}\theta \mathrm{d}\phi$$

- For axisymmetric radiation field,

$$F_{\lambda} = 2\pi \int_{0}^{\pi} I_{\lambda} \cos \theta \sin \theta d\theta = 2\pi \int_{-1}^{1} I_{\lambda} \mu d\mu$$

- Often, split flux for axisymmetric field into up ($\mu > 0$) and down ($\mu < 0$) components:

$$F_{\lambda} = F_{\lambda,+} - F_{\lambda,-}$$

where

$$F_{\lambda,+} = 2\pi \int_0^1 I_{\lambda} \mu d\mu,$$

$$F_{\lambda,-} = -2\pi \int_{-1}^0 I_{\lambda} \mu d\mu$$

- For isotropic radiation field,

$$F_{\lambda} = 2\pi I_{\lambda} \int_{-1}^{1} \mu \mathrm{d}\mu = 0$$

(and $F_{\lambda,+} = -F_{\lambda,-}$)

- Radiation pressure
 - Pressure can be thought of as amount of **momentum** in direction normal to surface, passing through unit area of surface each second
 - Photons carry momentum: p = E/c
 - Apply to radiation field: for energy E_{λ} passing through surface per second per unit area, momentum in normal direction is

$$p_{\lambda,\perp} = \frac{E_{\lambda}}{c}\cos\theta$$

- Integrate over all solid angles to get radiation pressure

$$P_{\mathrm{rad},\lambda} = \frac{\int E_{\lambda} \cos \theta}{c \mathrm{d}A \mathrm{d}t \mathrm{d}\lambda} = \frac{1}{c} \int I_{\lambda} \cos^2 \theta \mathrm{d}\Omega$$

- Note: the radiation pressure exists irrespective of whether the photons interact with matter. There is net force due to radiation when there is a gradient of radiation pressure (i.e., imbalance in radiation pressure on opposite sides of region).
- For axisymmetric radiation field,

$$P_{\mathrm{rad},\lambda} = \frac{2\pi}{c} \int_0^{\pi} I_\lambda \cos^2 \theta \sin \theta \mathrm{d}\theta = \frac{2\pi}{c} \int_{-1}^1 I_\lambda \mu^2 \mathrm{d}\mu$$

- For isotropic radiation field,

$$P_{\mathrm{rad},\lambda} = \frac{2\pi}{c} I_{\lambda} \int_{-1}^{1} \mu^2 \mathrm{d}\mu = \frac{4\pi}{3c} I_{\lambda}$$

- Total radiation pressure (integrated over all wavelengths):

$$P_{\rm rad} = \int_0^\infty P_{{\rm rad},\lambda} \mathrm{d}\lambda$$

For isotropic,

$$P_{\rm rad} = \frac{1}{3}u$$

(compare with monatomic gas: P = 2/3u)